

Lecture PowerPoints

Chapter 8

Physics: Principles with Applications, 6th edition

Giancoli

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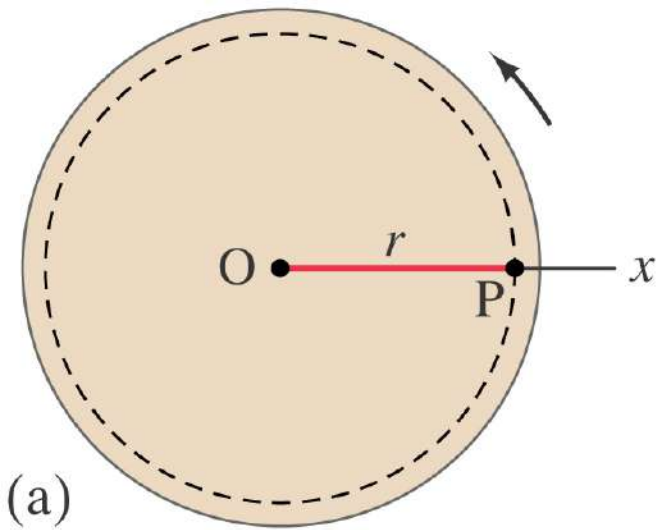
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Chapter 8

Rotational Motion

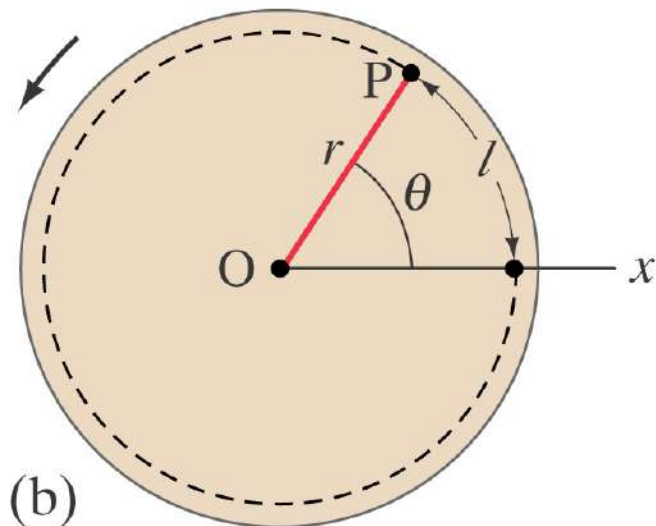


8-1 Angular Quantities



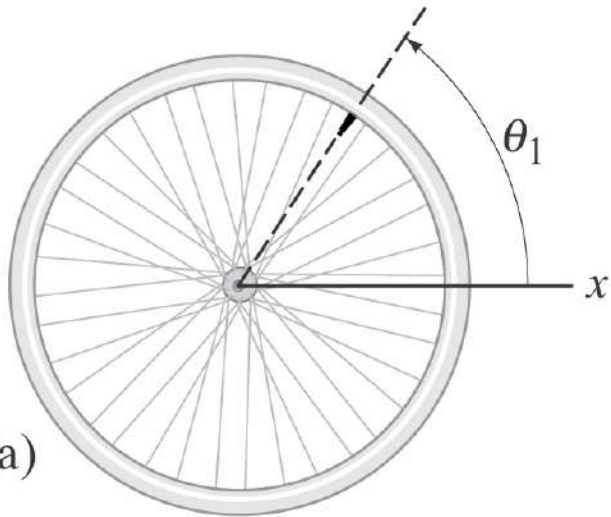
In purely rotational motion, all points on the object move in circles around the axis of rotation (“O”). The radius of the circle is r . All points on a straight line drawn through the axis move through the same angle in the same time. The angle θ in radians is defined:

$$\theta = \frac{l}{r} \quad (8-1a)$$



where l is the arc length.

8-1 Angular Quantities



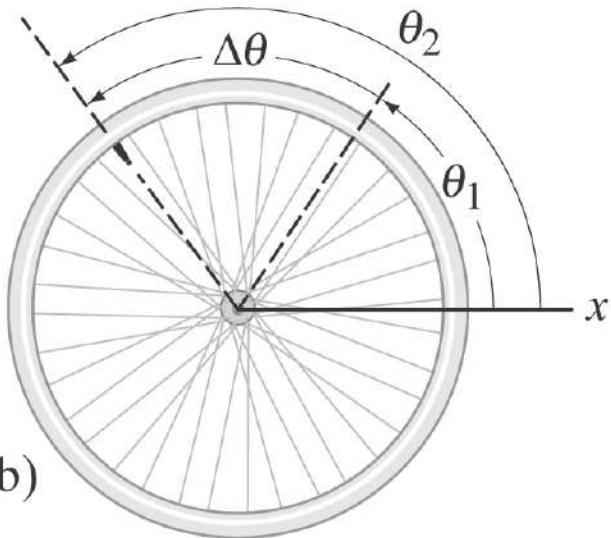
(a)

Angular displacement:

$$\Delta\theta = \theta_2 - \theta_1$$

The average angular velocity is defined as the total angular displacement divided by time:

$$\bar{\omega} = \frac{\Delta\theta}{\Delta t} \quad (8-2a)$$



(b)

The instantaneous angular velocity:

$$\omega = \lim_{\Delta t \rightarrow 0} \frac{\Delta\theta}{\Delta t} \quad (8-2b)$$

8-1 Angular Quantities

The **angular acceleration** is the rate at which the **angular velocity** changes with time:

$$\bar{\alpha} = \frac{\omega_2 - \omega_1}{\Delta t} = \frac{\Delta\omega}{\Delta t} \quad (8-3a)$$

The **instantaneous acceleration**:

$$\alpha = \lim_{\Delta t \rightarrow 0} \frac{\Delta\omega}{\Delta t} \quad (8-3b)$$

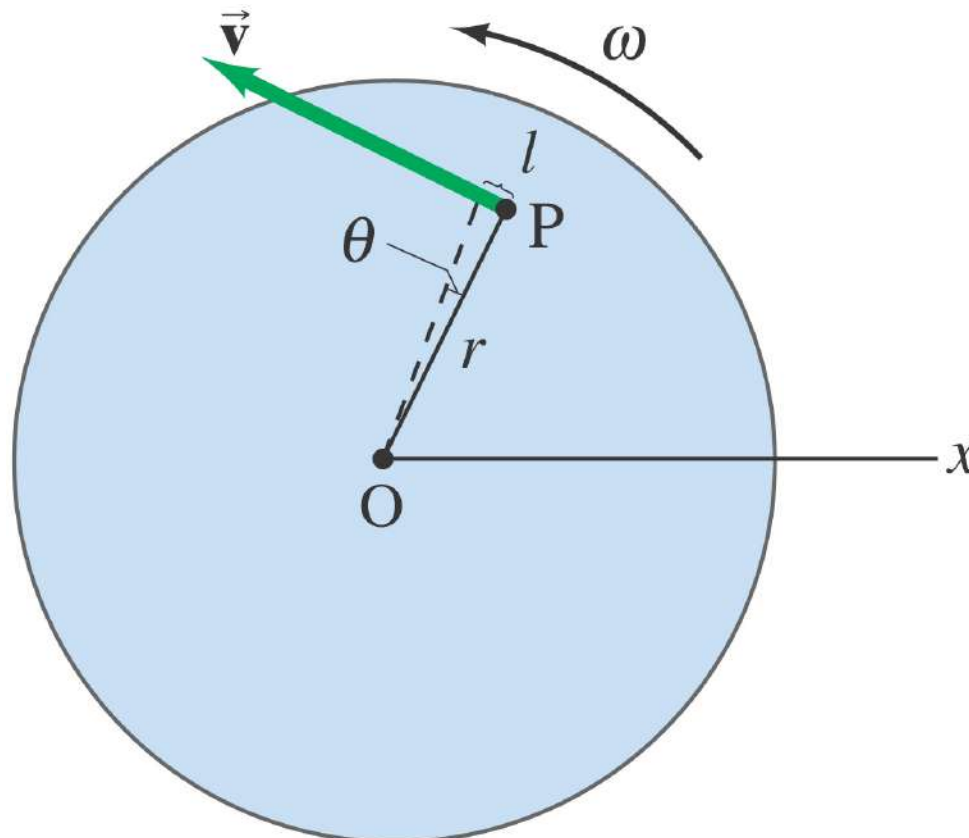
8-1 Angular Quantities

Every point on a rotating body has an angular velocity ω and a linear velocity v .

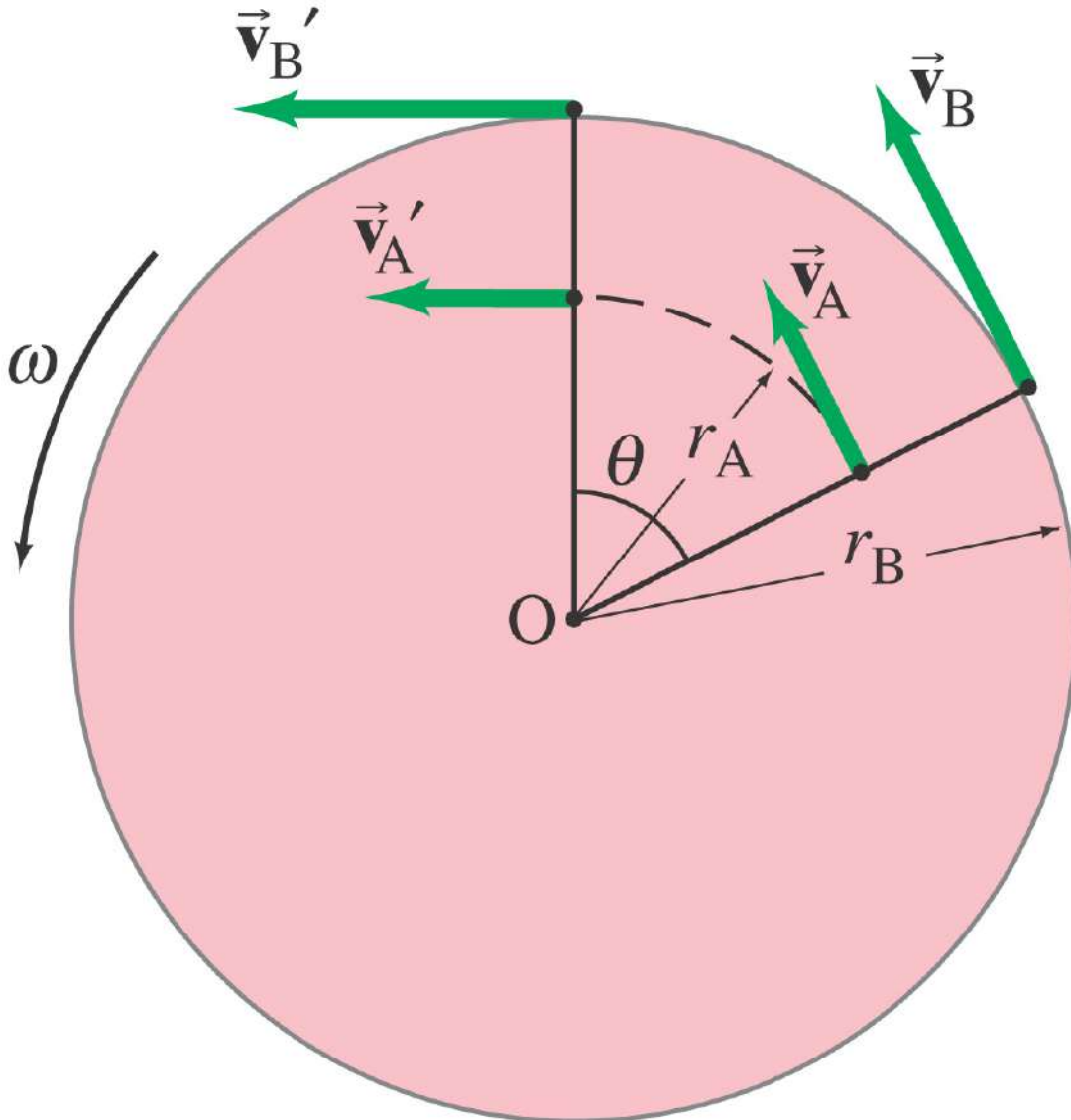
They are related:

$$v = r\omega$$

(8-4)

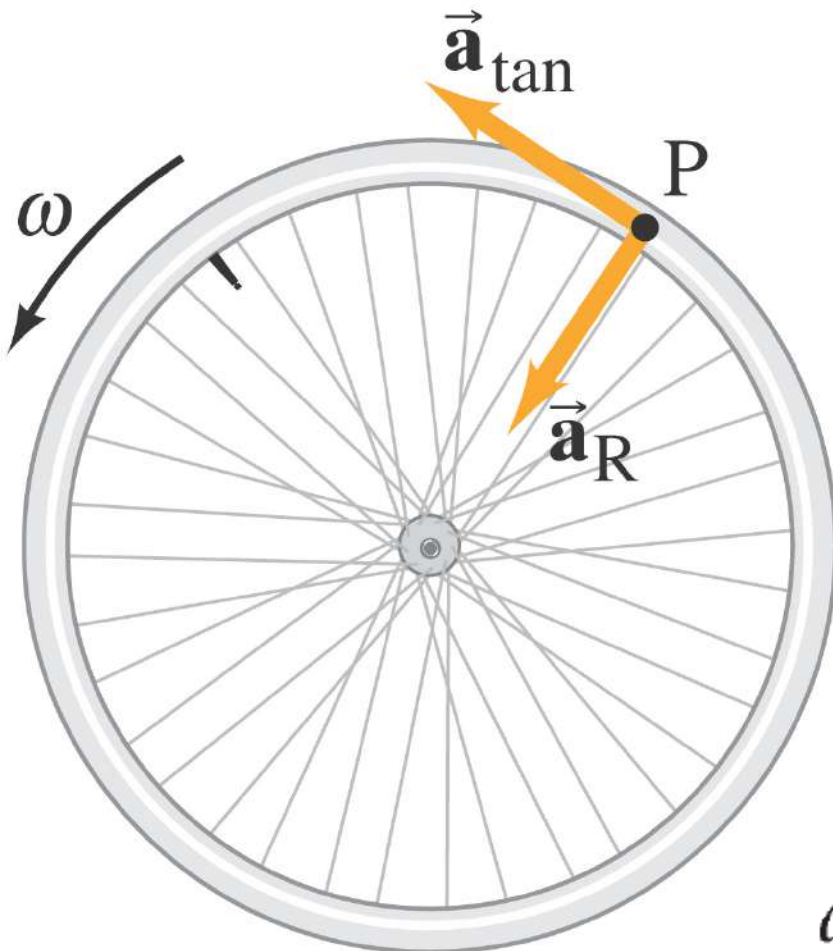


8-1 Angular Quantities



Therefore, objects farther from the axis of rotation will move faster.

8-1 Angular Quantities



If the angular velocity of a rotating object changes, it has a **tangential acceleration**:

$$a_{\text{tan}} = r\alpha \quad (8-5)$$

Even if the angular velocity is constant, each point on the object has a **centripetal acceleration**:

$$a_R = \frac{v^2}{r} = \frac{(r\omega)^2}{r} = \omega^2 r \quad (8-6)$$

8-1 Angular Quantities

Here is the correspondence between linear and rotational quantities:

TABLE 8–1 Linear and Rotational Quantities

Linear	Type	Rotational	Relation
x	displacement	θ	$x = r\theta$
v	velocity	ω	$v = r\omega$
a_{tan}	acceleration	α	$a_{\text{tan}} = r\alpha$

8-1 Angular Quantities

The frequency is the number of complete revolutions per second: $f = \frac{\omega}{2\pi}$

Frequencies are measured in hertz.

$$1 \text{ Hz} = 1 \text{ s}^{-1}$$

The period is the time one revolution takes:

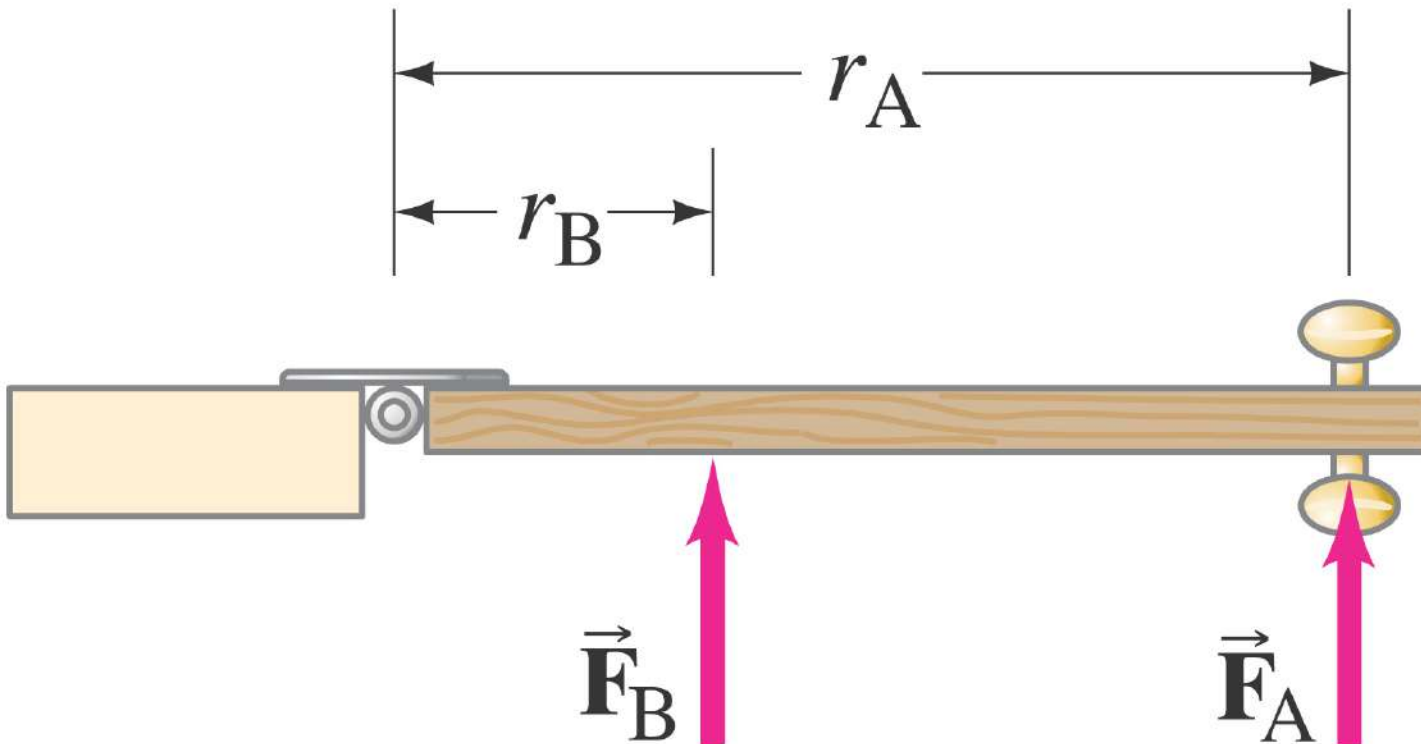
$$T = \frac{1}{f}$$

(8-8)

8-4 Torque

To make an object start rotating, a force is needed; the position and direction of the force matter as well.

The perpendicular distance from the axis of rotation to the line along which the force acts is called the lever arm.



8-4 Torque



(a)

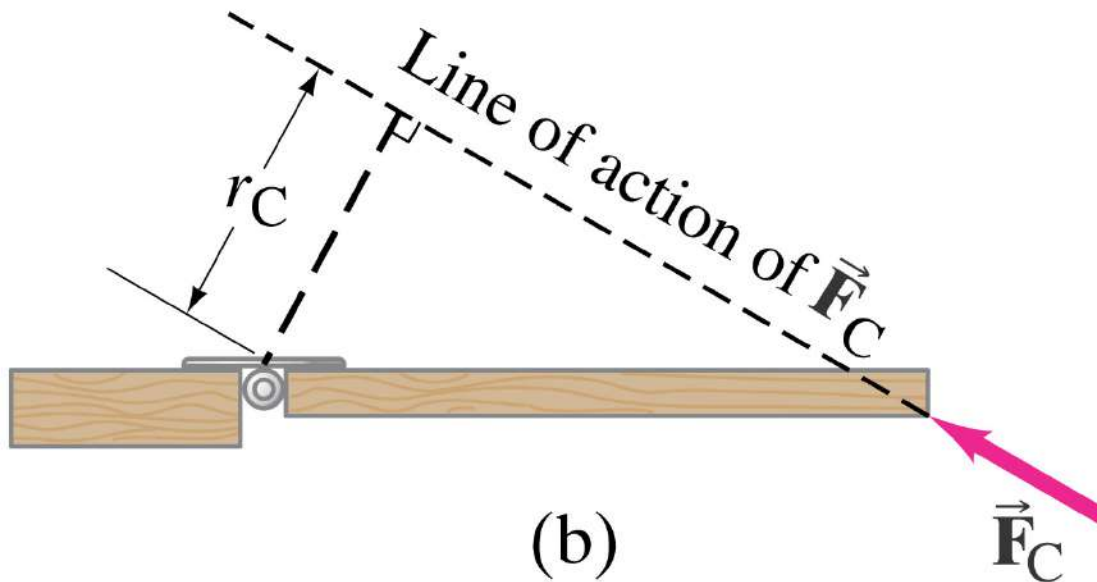
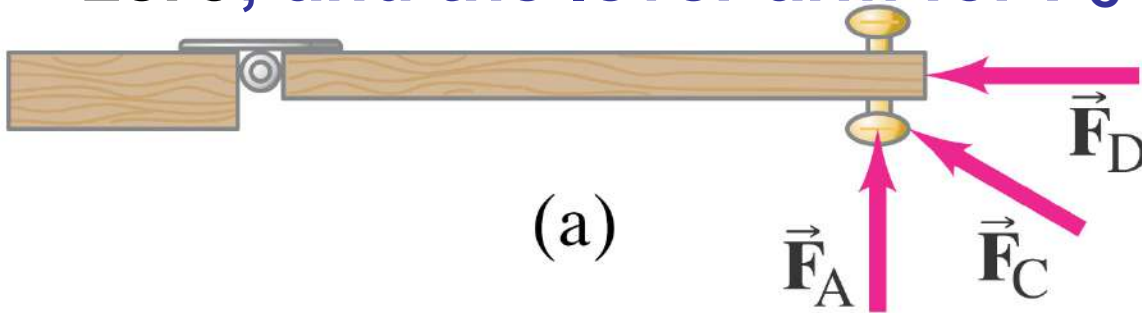


(b)

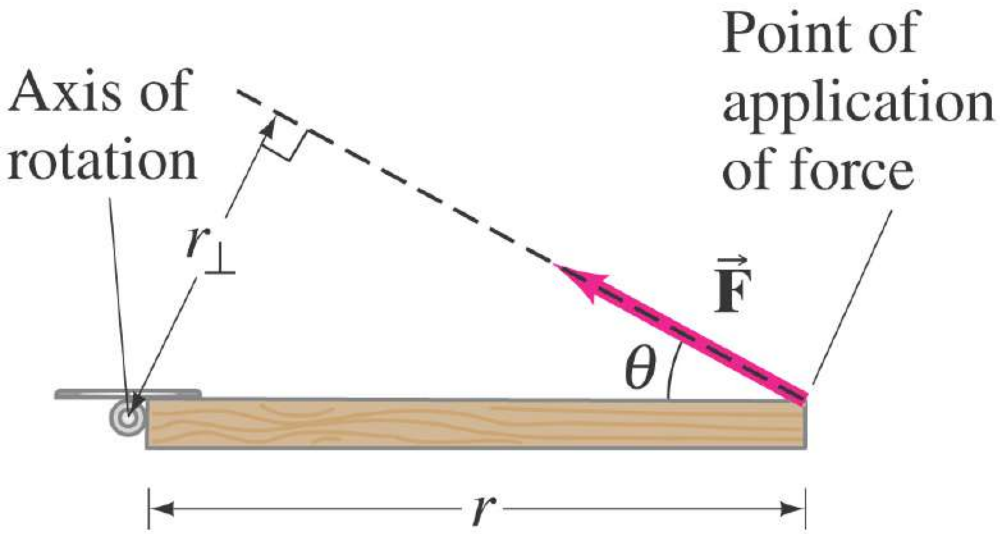
A longer lever arm is very helpful in rotating objects.

8-4 Torque

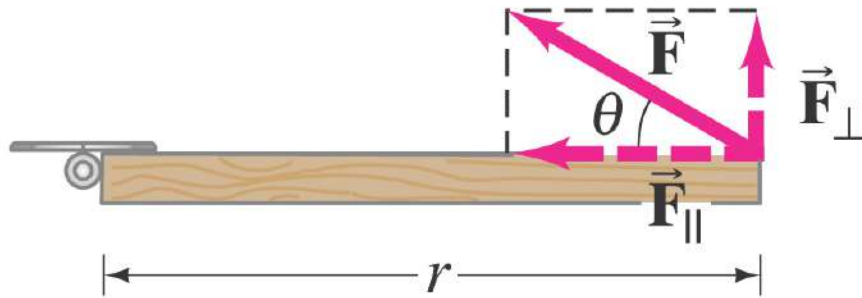
Here, the lever arm for F_A is the distance from the knob to the hinge; the lever arm for F_D is zero; and the lever arm for F_C is as shown.



8-4 Torque



(a)



(b)

The torque is defined as:

$$\tau = r_{\perp} F$$

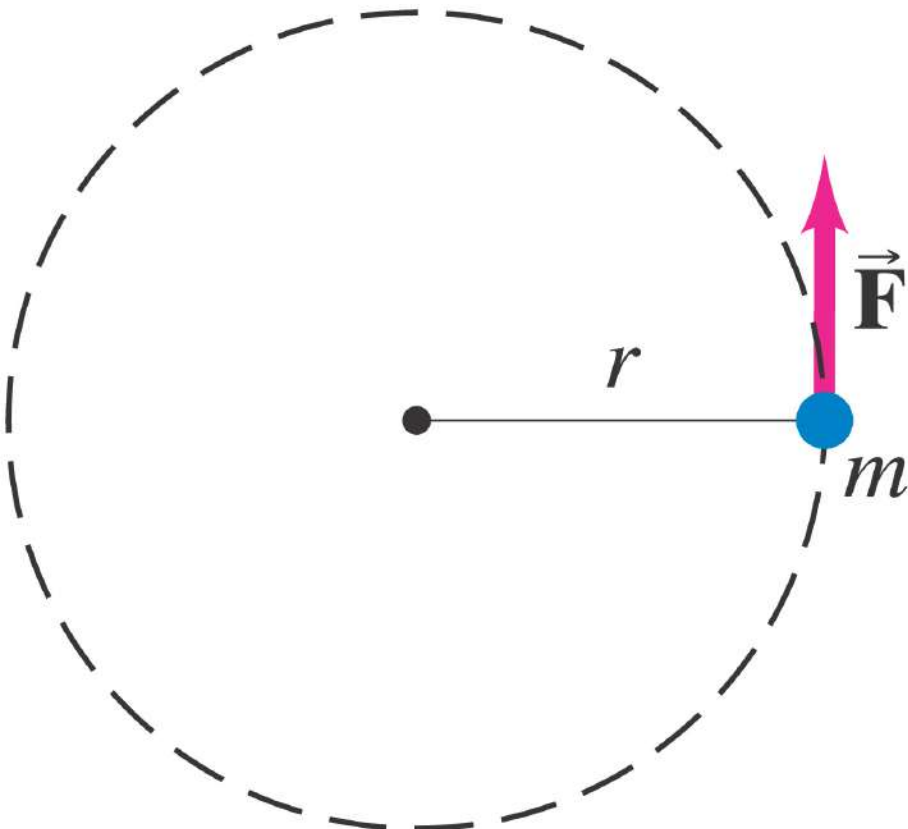
$$= r F \sin \theta$$

(on formula sheet)

8-5 Rotational Dynamics; Torque and Rotational Inertia

Knowing that $F = ma$
and $a_{\text{tan}} = r\alpha$

$$\tau = mr^2\alpha \quad (8-11)$$



This is for a single point mass; what about an extended object?

As the angular acceleration is the same for the whole object, we can write:

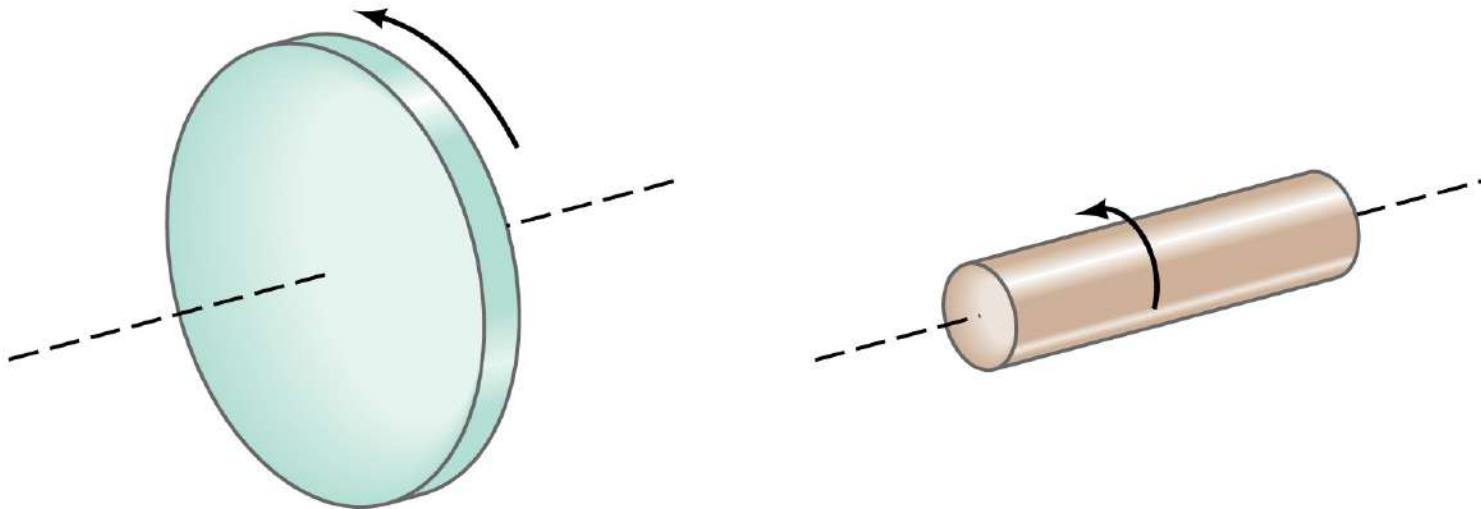
$$\Sigma\tau = (\Sigma mr^2)\alpha \quad (8-12)$$

8-5 Rotational Dynamics; Torque and Rotational Inertia

$$I = \sum mr^2$$

The quantity is called the **rotational inertia** (or **moment of inertia**) of an object. **Moment of Inertia** is the resistance to angular acceleration.

The **distribution of mass** matters here – these two objects have the same mass, but the one on the left has a **greater rotational inertia**, as so much of its mass is far from the axis of rotation.



8-5 Rotational Dynamics; Torque and Rotational Inertia

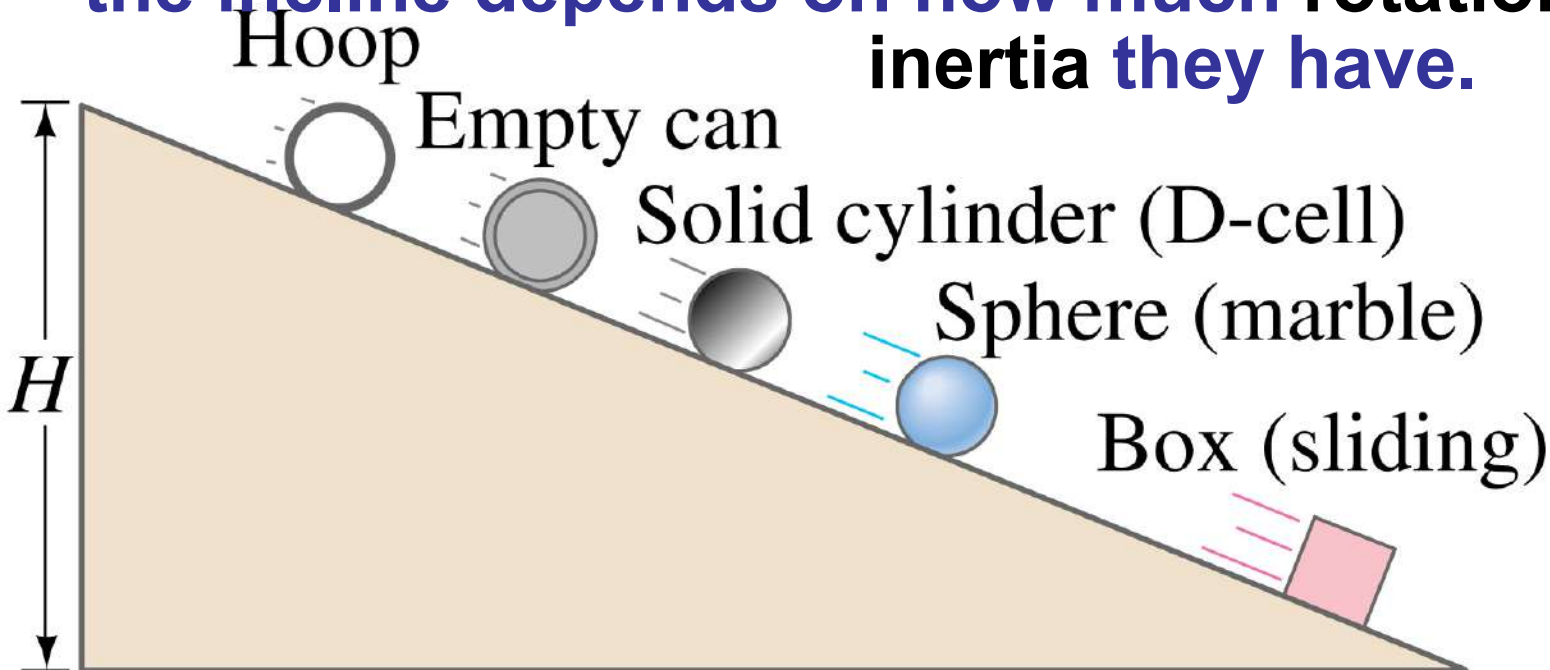
Object	Location of axis	Moment of inertia
(a) Thin hoop, radius R	Through center	MR^2
(b) Thin hoop, radius R width W	Through central diameter	$\frac{1}{2}MR^2 + \frac{1}{12}MW^2$
(c) Solid cylinder, radius R	Through center	$\frac{1}{2}MR^2$
(d) Hollow cylinder, inner radius R_1 outer radius R_2	Through center	$\frac{1}{2}M(R_1^2 + R_2^2)$
(e) Uniform sphere, radius R	Through center	$\frac{2}{5}MR^2$
(f) Long uniform rod, length L	Through center	$\frac{1}{12}ML^2$
(g) Long uniform rod, length L	Through end	$\frac{1}{3}ML^2$
(h) Rectangular thin plate, length L , width W	Through center	$\frac{1}{12}M(L^2 + W^2)$

The rotational inertia of an object depends not only on its mass distribution but also the location of the axis of rotation – compare (f) and (g), for example.

8-7 Rotational Kinetic Energy

When using **conservation of energy**, both **rotational and translational kinetic energy** must be taken into account.

All these objects have the same **potential energy** at the top, but the time it takes them to get down the **incline** depends on how much **rotational inertia** they have.



8-8 Angular Momentum and Its Conservation

In analogy with linear momentum, we can define angular momentum L :

$$L = I\omega \quad (8-18)$$

We can then write the total torque as being the rate of change of angular momentum.

If the net torque on an object is zero, the total angular momentum is constant.

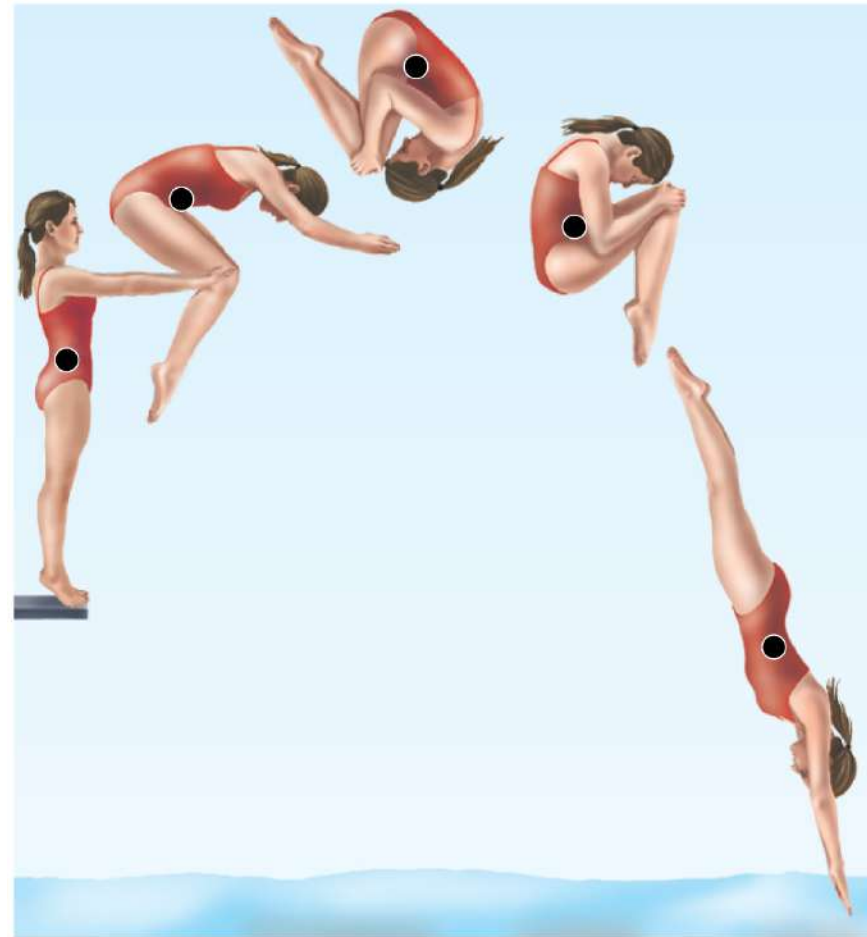
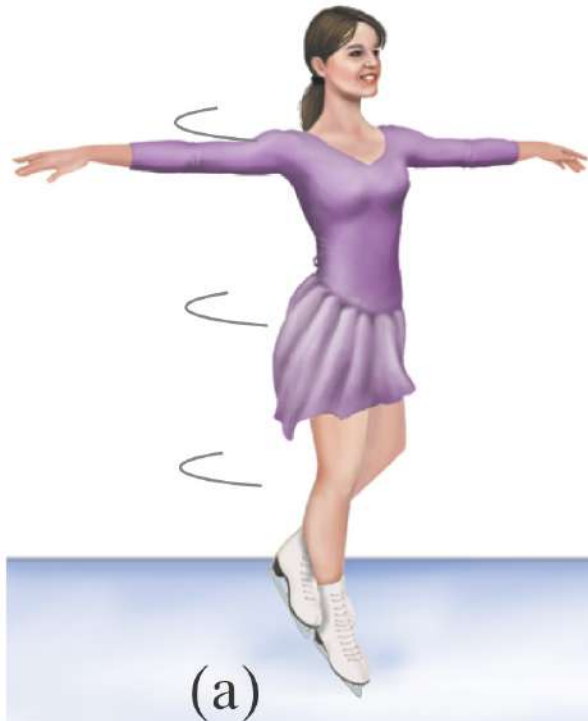
$$I\omega = I_0\omega_0 = \text{constant}$$

8-8 Angular Momentum and Its Conservation

Therefore, systems that can change their rotational inertia through internal forces will also change their rate of rotation:

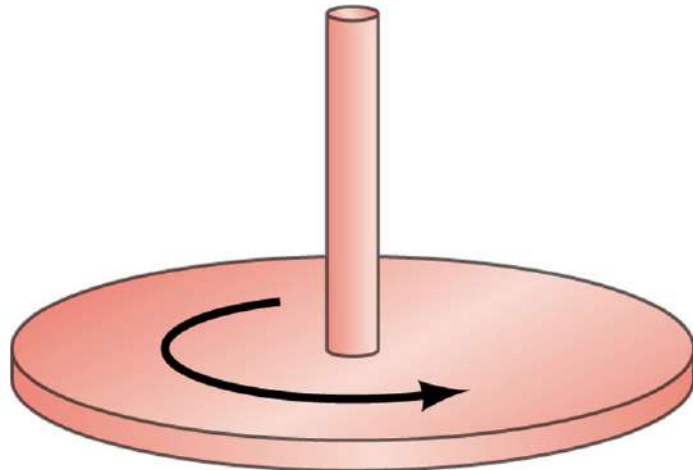
I large,
 ω small

I small,
 ω large

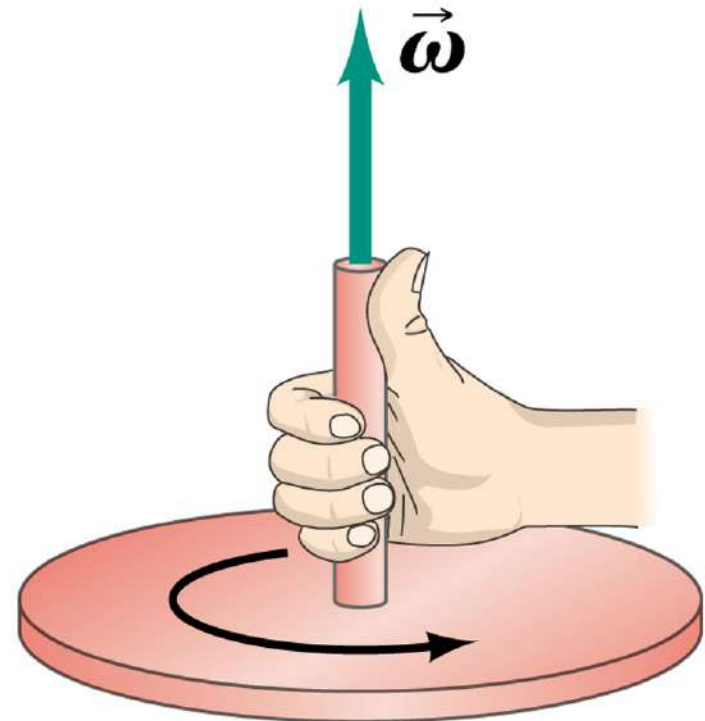


8-9 Vector Nature of Angular Quantities

The angular velocity vector points along the axis of rotation; its direction is found using a right hand rule:



(a)



(b)